A numerical approach to stochastic inflation and primordial black holes

TAUP2021

Eemeli Tomberg, NICPB Tallinn eemeli.tomberg@kbfi.ee

Based on 2012.06551, in collaboration with D. Figueroa, S. Raatikainen, S. Räsänen



Cosmic inflation



Cosmic inflation

Cosmological perturbations

Concepts

Cosmic inflation

Cosmological perturbations

Primordial black holes



Stochastic inflation



Stochastic inflation

Includes non-linear effects

Concepts

Stochastic inflation

- Includes non-linear effects
- Numerical method: even more non-linearities

Model of inflation fits CMB



Model of inflation fits CMB



Model of inflation fits CMB

Origin of perturbations: fluctuations of quantum vacuum

Origin of perturbations: fluctuations of quantum vacuum

Space expands and perturbations get stretched

Origin of perturbations: fluctuations of quantum vacuum

Space expands and perturbations get stretched

Perturbations (eventually) become classical and freeze after crossing Hubble horizon

Origin of perturbations: fluctuations of quantum vacuum

Space expands and perturbations get stretched

Perturbations (eventually) become classical and freeze after crossing Hubble horizon

Strong perturbations from ultra-slow-roll inflation

Must go beyond linear perturbations

Coarse-grain perturbations over super-Hubble scales

Gradient expansion: to leading order, coarse-grained perturbations follow locally (non-linear) FLRW equations [Class.Q.Grav.9,1943(1992)]

Must go beyond linear perturbations

Coarse-grain perturbations over super-Hubble scales

Gradient expansion: to leading order, coarse-grained perturbations follow locally (non-linear) FLRW equations [Class.Q.Grav.9,1943(1992)]

 ΔN formalism: from FLRW variables to perturbation variables [astro-ph/9507001]

• Change in e-folds of expansion $\Delta N =$ curvature perturbation \mathcal{R}

Stretching perturbations give stochastic kicks

When perturbations of a certain scale stretch to the coarse-graining scale, they get coarse-grained

Stretching perturbations give stochastic kicks

When perturbations of a certain scale stretch to the coarse-graining scale, they get coarse-grained

Result: 'kicks' to coarse-grained field. Random due to quantum initial conditions

Stretching perturbations give stochastic kicks

When perturbations of a certain scale stretch to the coarse-graining scale, they get coarse-grained

Result: 'kicks' to coarse-grained field. Random due to quantum initial conditions

Stochastic evolution of local coarse-grained field [Lect.Notes Phys.246,107(1986)]

PBHs form from strong perturbations

During radiation domination, perturbations re-enter Hubble radius

PBHs form from strong perturbations

During radiation domination, perturbations re-enter Hubble radius

Perturbation collapses to black hole if it exceeds threshold [1309.4201, 1405.7023, 2011.03014]

PBHs form from strong perturbations

During radiation domination, perturbations re-enter Hubble radius

Perturbation collapses to black hole if it exceeds threshold [1309.4201, 1405.7023, 2011.03014]

BH mass = all the mass inside one Hubble radius when the scale re-enters

 Initial conditions: CMB scale, Bunch–Davies vacuum

 Initial conditions: CMB scale, Bunch–Davies vacuum

Stochastic evolution with backreaction

Non-linear interactions included

Non-linear interactions included

Compare to simpler approach with noise $\sim \frac{H^2}{2\pi^2}$

 Initial conditions: CMB scale, Bunch–Davies vacuum

Stochastic evolution with backreaction

 Initial conditions: CMB scale, Bunch–Davies vacuum

Stochastic evolution with backreaction

Stochastic kicks end when PBH scale reached

 Initial conditions: CMB scale, Bunch–Davies vacuum

Stochastic evolution with backreaction

Stochastic kicks end when PBH scale reached

Continue (without kicks) to constant- ϕ hypersurface, record N

 Initial conditions: CMB scale, Bunch–Davies vacuum

Stochastic evolution with backreaction

Stochastic kicks end when PBH scale reached

Continue (without kicks) to constant- ϕ hypersurface, record N

Repeat 10^{11} times, collect statistics

Scale $M_{\rm PBH} = 10^{-14} M_{\odot}$, $k_{\rm PBH} = 10^{13} \rm ~Mpc^{-1}$ chosen so that USR ends when $k_{\rm PBH}$ gets coarse-grained

Scale $M_{\rm PBH} = 10^{-14} M_{\odot}$, $k_{\rm PBH} = 10^{13} \rm \ Mpc^{-1}$ chosen so that USR ends when $k_{\rm PBH}$ gets coarse-grained

To contribute significantly to dark matter, need initial fraction $\beta \sim 10^{-16}$

Scale $M_{\rm PBH} = 10^{-14} M_{\odot}$, $k_{\rm PBH} = 10^{13} \rm \ Mpc^{-1}$ chosen so that USR ends when $k_{\rm PBH}$ gets coarse-grained

To contribute significantly to dark matter, need initial fraction $\beta \sim 10^{-16}$

Gaussian statistics:

$$\sigma_{\mathcal{R}}^{2} = \int^{k_{\text{PBH}}} \mathrm{d}(\ln k) \mathcal{P}_{\mathcal{R}}(k)$$
$$\beta = 2 \int_{\mathcal{R}_{c}}^{\infty} \mathrm{d}\mathcal{R} \frac{1}{\sqrt{2\pi}\sigma_{\mathcal{R}}} e^{-\frac{\mathcal{R}^{2}}{2\sigma_{\mathcal{R}}^{2}}} \approx \frac{\sqrt{2}\sigma_{\mathcal{R}}}{\sqrt{\pi}\mathcal{R}_{c}} e^{-\frac{\mathcal{R}^{2}}{2\sigma_{\mathcal{R}}^{2}}}$$

True abundance much higher than Gaussian estimate

Numerics: exponential tail, with $\beta = 1.2 \times 10^{-10}, \quad \Omega_{\rm PBH} = 5.4 \times 10^4$

Larger than Gaussian result by factor $10^5!$

Future directions

More statistics

More models

Full mass spectrum

Correlations between different scales

Conclusions

Stochastic inflation captures non-linearities of cosmological perturbations

Crucial for PBH formation

Introduced a numerical recipe to calculate these in a general single-field model

Thank you!

[2012.06551]