

Preheating in Palatini Higgs inflation

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[1902.10148]

In collaboration with Javier Rubio

Contents

- ▶ Higgs inflation: metric vs Palatini formalisms
- ▶ Reheating
- ▶ Significance on cosmology

Motivation

- ▶ Planck: spectral index $n_s = 0.9625 \pm 0.0048$
[1807.06211]
- ▶ Typical inflationary models: $n_s \approx 1 - aN^b$,
 $N \sim 50$ is number of e-folds of inflation

$$\Delta n_s \approx -aN^{b-1} \Delta N \approx b(n_s - 1) \frac{\Delta N}{N} \sim 10^{-3} \Delta N$$

General relativity: degrees of freedom

- ▶ Metric $g_{\mu\nu}$ gives local lengths and angles
- ▶ Connection $\Gamma_{\beta\gamma}^{\alpha}$ gives covariant derivatives
- ▶ Metric formulation: Levi-Civita connection

$$\Gamma_{\beta\gamma}^{\alpha} = \frac{1}{2}g^{\alpha\rho}(g_{\beta\rho,\gamma} + g_{\gamma\rho,\beta} - g_{\beta\gamma,\rho})$$

- ▶ Palatini formulation: $g_{\mu\nu}$ and $\Gamma_{\beta\gamma}^{\alpha}$ independent

Higgs inflation

- ▶ Non-minimal coupling to gravity

$$S = \int d^4x \sqrt{-g} \left[-\frac{1}{2} (M^2 + \xi h^2) \underbrace{g^{\mu\nu}}_{g_{\mu\nu}} R_{\mu\nu} + \frac{1}{2} \underbrace{g^{\mu\nu}}_{g_{\mu\nu}} \partial_\mu h \partial_\nu h - \frac{\lambda}{4} h^4 \right]$$

$\Gamma_{\beta\gamma}^\alpha$

Higgs inflation: metric formulation

- ▶ Non-minimal coupling to gravity

$$S = \int d^4x \sqrt{-g} \left[\underbrace{-\frac{1}{2}}_{g_{\mu\nu}} (M^2 + \xi h^2) \underbrace{g^{\mu\nu}}_{g_{\mu\nu}} R_{\mu\nu} + \frac{1}{2} \underbrace{g^{\mu\nu}}_{g_{\mu\nu}} \partial_\mu h \partial_\nu h - \frac{\lambda}{4} h^4 \right]$$

- ▶ Weyl transformation to Einstein frame:

$$g_{\mu\nu} = g_{E\mu\nu} \left(1 + \frac{\xi h^2}{M^2} \right)^{-2}$$
$$R_{\mu\nu} = R_{E\mu\nu} + \dots$$

Higgs inflation: metric formulation

- ▶ Einstein frame action:

$$S = \int d^4x \sqrt{-g_E} \left[-\frac{1}{2} M^2 g^{E\mu\nu} R_{E\mu\nu} + \frac{1}{2} g^{E\mu\nu} \frac{1 + \xi h^2/M^2 + 6\xi^2 h^2/M^2}{(1 + \xi h^2/M^2)^2} \partial_\mu h \partial_\nu h - \frac{\lambda h^4}{4 \left(1 + \frac{\xi h^2}{M^2}\right)^2} \right]$$

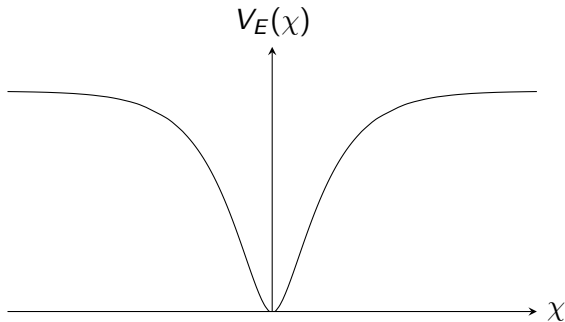
$\frac{d\chi}{dh} = \frac{\sqrt{1 + \xi h^2 + 6\xi^2 h^2}}{1 + \xi h^2}$ $V_E(h)$

Higgs inflation: metric formulation

- ▶ Einstein frame action:

$$S = \int d^4x \sqrt{-g_E} \left[-\frac{1}{2} M^2 g^{E\mu\nu} R_{E\mu\nu} + \frac{1}{2} g^{E\mu\nu} \partial_\mu \chi \partial_\nu \chi - V_E(\chi) \right]$$

- ▶ Einstein frame potential:



Higgs inflation: Palatini formulation

- ▶ Non-minimal coupling to gravity

$$S = \int d^4x \sqrt{-g} \left[-\frac{1}{2} (M^2 + \xi h^2) \underbrace{g^{\mu\nu}}_{g_{\mu\nu}} R_{\mu\nu} + \frac{1}{2} \underbrace{g^{\mu\nu}}_{g_{\mu\nu}} \partial_\mu h \partial_\nu h - \frac{\lambda}{4} h^4 \right]$$

$\Gamma_{\beta\gamma}^\alpha$

- ▶ Weyl transformation to Einstein frame:

$$g_{\mu\nu} = g_{E\mu\nu} \left(1 + \frac{\xi h^2}{M^2} \right)^{-2}$$
$$R_{\mu\nu} = R_{E\mu\nu}$$

Higgs inflation: Palatini formulation

- ▶ Einstein frame action:

$$S = \int d^4x \sqrt{-g_E} \left[-\frac{1}{2} M^2 g^{E\mu\nu} R_{E\mu\nu} + \right. \\ \left. \frac{1}{2} g^{E\mu\nu} \frac{1}{1+\xi h^2/M^2} \partial_\mu h \partial_\nu h - \frac{\lambda h^4}{4 \left(1 + \frac{\xi h^2}{M^2}\right)^2} \right]$$

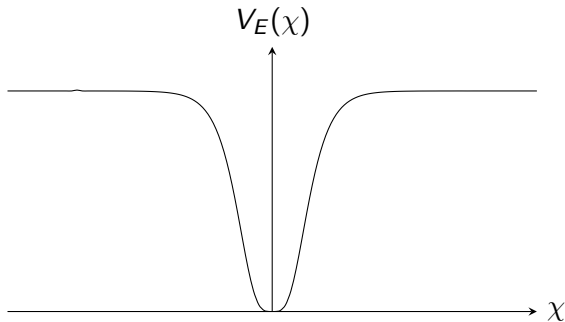
$\frac{d\chi}{dh} = \frac{1}{\sqrt{1+\xi h^2}}$ $V_E(h)$

Higgs inflation: Palatini formulation

- ▶ Einstein frame action:

$$S = \int d^4x \sqrt{-g_E} \left[-\frac{1}{2} M^2 g^{E\mu\nu} R_{E\mu\nu} + \frac{1}{2} g^{E\mu\nu} \partial_\mu \chi \partial_\nu \chi - V_E(\chi) \right]$$

- ▶ Einstein frame potential:



Metric vs. Palatini

- ▶ Metric: [0710.3755]

$$\xi \approx 800\sqrt{\lambda}N \sim 10^4 \quad (N = 50, \lambda \sim 0.1)$$

$$n_s \approx 1 - \frac{2}{N} \approx 0.96, \quad r \approx \frac{12}{N^2} \approx 4.8 \times 10^{-3}$$

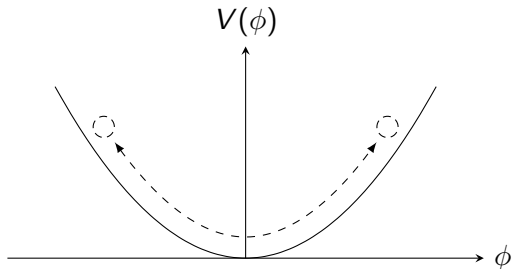
- ▶ Palatini: [0803.2664]

$$\xi \approx 3.8 \times 10^6 \lambda N^2 \sim 10^9$$

$$n_s \approx 1 - \frac{2}{N} \approx 0.96, \quad r \approx \frac{2}{\xi N^2} \sim 10^{-12}$$

Reheating: Overview

- ▶ After inflation, inflaton oscillates around its minimum



$$\ddot{\chi} + \underbrace{3H\dot{\chi}}_{\text{friction}} + \underbrace{V'(\chi)}_{\text{oscillation}} = 0$$

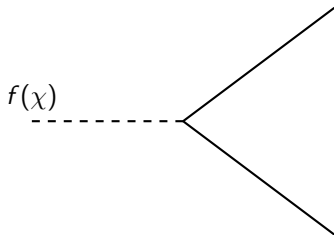
friction oscillation

Reheating: Particle production

- ▶ Need to transfer some of the energy density in the inflaton condensate to particles
- ▶ This happens through interactions of the inflaton field

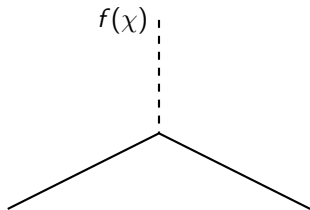
Higgs reheating

- ▶ Interactions known
- ▶ Decays of Higgs condensate (subdominant)



Higgs reheating

- ▶ Higgs-induced mass terms for fermions, weak gauge bosons and Higgs bosons



$$m_W^2 = \frac{g^2 h^2}{4(1 + \xi h^2/M^2)}, \quad m_Z^2 = \frac{(g^2 + g'^2) h^2}{4(1 + \xi h^2/M^2)},$$
$$m_t^2 = \frac{y_t^2 h^2}{2(1 + \xi h^2/M^2)}, \quad m_h^2 = V_E''(h)$$

Time-dependent mass: Preheating

- ▶ Scalar field mode function:

$$\ddot{Q}_k + 3H\dot{Q}_k + \underbrace{\left[\frac{k^2}{a^2} + m^2(\chi) \right]}_{\omega_k^2} Q_k = 0$$

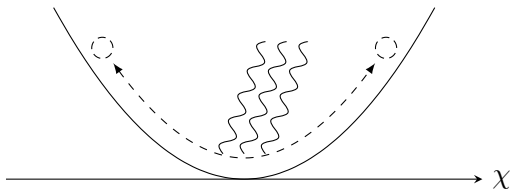
- ▶ For constant m , standard solution:

$$a^{3/2} Q_k = \frac{\alpha_k}{\sqrt{2\omega_k}} e^{-i \int^t \omega_k(t') dt'} + \frac{\beta_k}{\sqrt{2\omega_k}} e^{i \int^t \omega_k(t') dt'}$$

- ▶ Adiabatic vacuum: $\alpha_k = 1, \beta_k = 0$
- ▶ Non-vacuum states: $n_k = |\beta_k|^2$

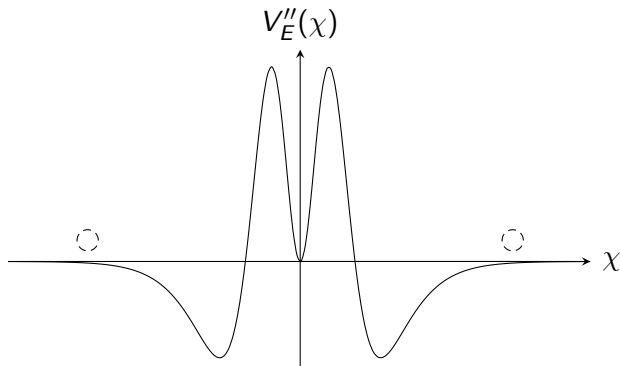
Parametric resonance

- ▶ Adiabaticity condition $|\dot{\omega}_k|/\omega_k^2 \gg 1$ broken near the bottom of the Higgs potential
- ▶ For certain Fourier modes: explosive particle production



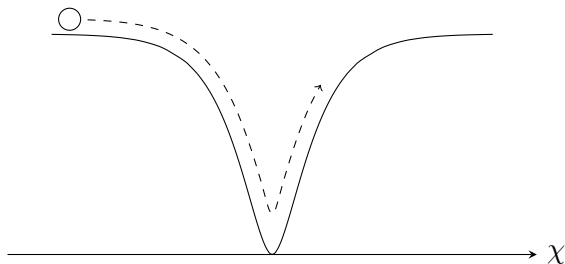
Tachyonic preheating

- ▶ If $m^2 < 0 \Rightarrow \omega_k^2 < 0$, mode function grows exponentially: $Q_k \propto e^{\sqrt{-\omega_k^2}t}$
- ▶ Possible for Higgs perturbations with $m_h^2 = V_E''$



Higgs preheating: Metric case

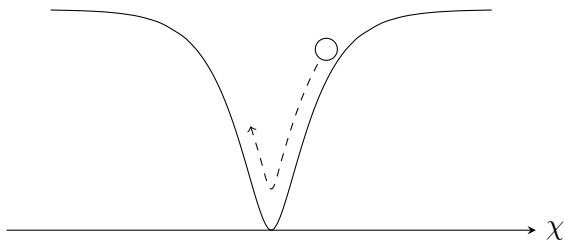
$$\ddot{\chi} + 3H\dot{\chi} + V'(\chi) = 0$$



- ▶ Quick decay of χ oscillation amplitude; then,
 $V_E = M\chi^2$, $M = \text{constant}$

Higgs preheating: Metric case

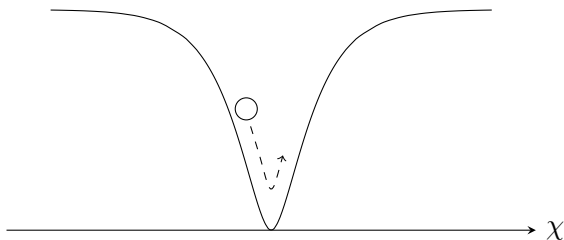
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Higgs preheating: Metric case

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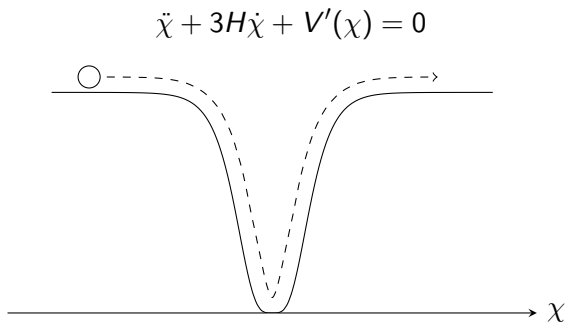


- ▶ Quick decay of χ oscillation amplitude; then,
 $V_E = M\chi^2$, $M=\text{constant}$

Higgs preheating: Metric case

- ▶ Higgs mass squared constant and positive: no production
- ▶ Parametric resonance: production of W , Z bosons
 - ▶ At early times, decay to fermions (*Combined reheating*)
 - ▶ Later, bosons start to accumulate
- ▶ Reheating takes a few e-folds

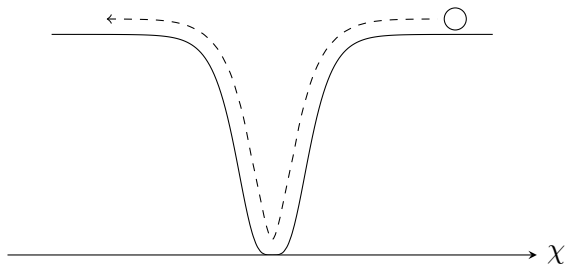
Higgs preheating: Palatini case



- ▶ Oscillation amplitude stays almost constant!
Why?

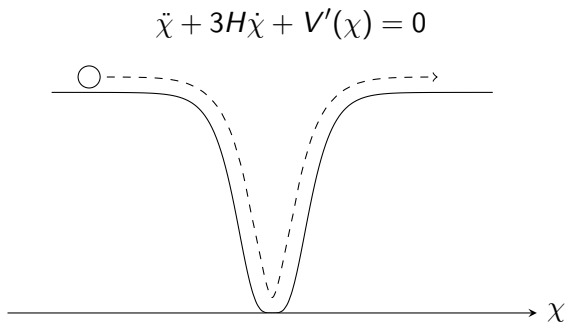
Higgs preheating: Palatini case

$$\ddot{\chi} + 3H\dot{\chi} + V'(\chi) = 0$$



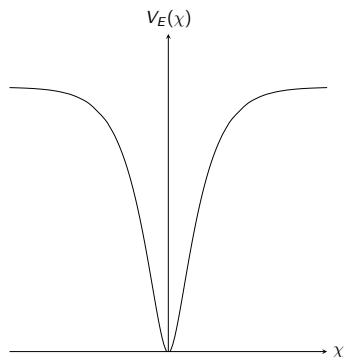
- ▶ Oscillation amplitude stays almost constant!
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Higgs preheating: Palatini case

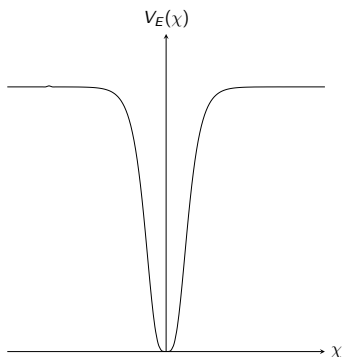


- ▶ Oscillation amplitude stays almost constant!
Why?

Comparison: Metric vs Palatini



$$\frac{d\chi}{dh} = \frac{\sqrt{1 + \xi h^2 + 6\xi^2 h^2}}{1 + \xi h^2}$$



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Comparison: Metric vs Palatini

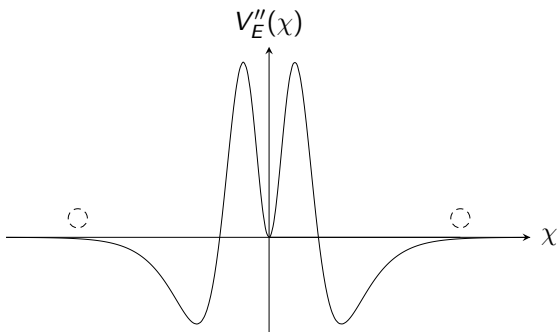
- ▶ Oscillation energy depleted more quickly in metric formulation:

$$M_P^2 \frac{dH}{dh} = -\frac{\dot{\chi}^2}{2\dot{h}} = -\frac{1}{2} \sqrt{6H^2 M_P^2 - 2V_E[\chi(h)]} \frac{d\chi}{dh}$$

Bigger in metric case

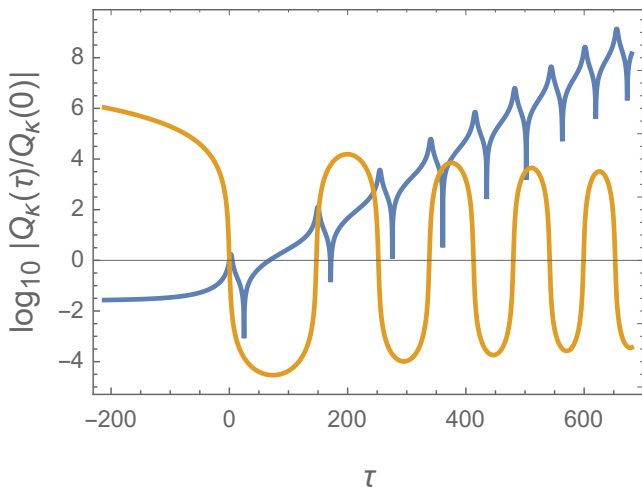
Higgs preheating: Palatini case

- ▶ Result: m_h^2 time-dependent, negative at times
- ▶ Tachyonic Higgs production is the leading preheating channel



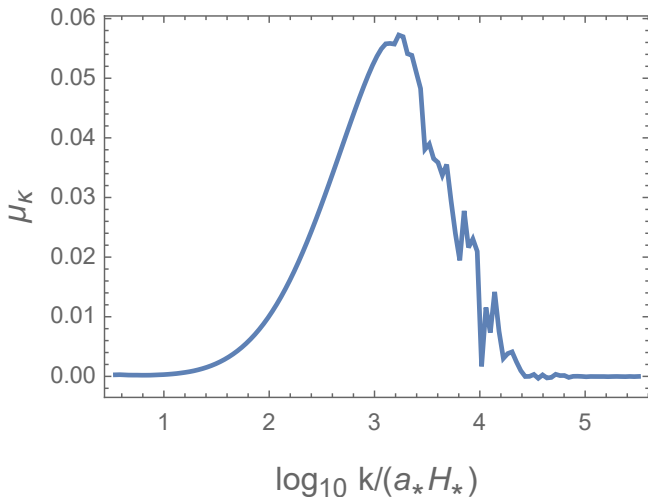
Higgs preheating: Palatini case

- ▶ Typical tachyonic mode function: $Q_k \propto e^{\mu_k \tau}$



Higgs preheating: Palatini case

- ▶ Growth index depends on k :



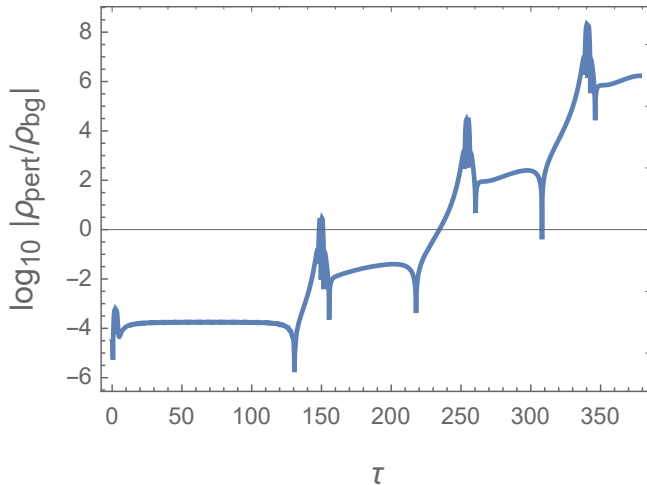
Higgs preheating: Palatini case

- ▶ Particle concept is ill-defined, but we can calculate energy density in perturbations:

$$\frac{\rho_{\text{pert}}}{\rho_{\text{B}}} = \frac{1}{3H^2 M_{\text{P}}^2} \int_{k_{\text{min}}}^{k_{\text{max}}} \frac{dk}{(2\pi)^3} \frac{4\pi k^2}{2} \left[|\dot{Q}_k|^2 + \left(\frac{k^2}{a^2} + V_E''(\chi) \right) |Q_k|^2 \right]$$

- ▶ Numerical calculations: significant fraction of energy density in perturbations after only a few oscillations, less than one e-fold

Higgs preheating: Palatini case



Significance on cosmology

- ▶ Palatini formulation: almost instant reheating

- ▶ Reheating temperature $T_{\text{RH}} = \left(\frac{30 \lambda}{4\pi^2 \xi^2 g_{* \text{RH}}} \right)^{1/4} M_P$

- ▶ E-folds of inflation from CMB pivot scale

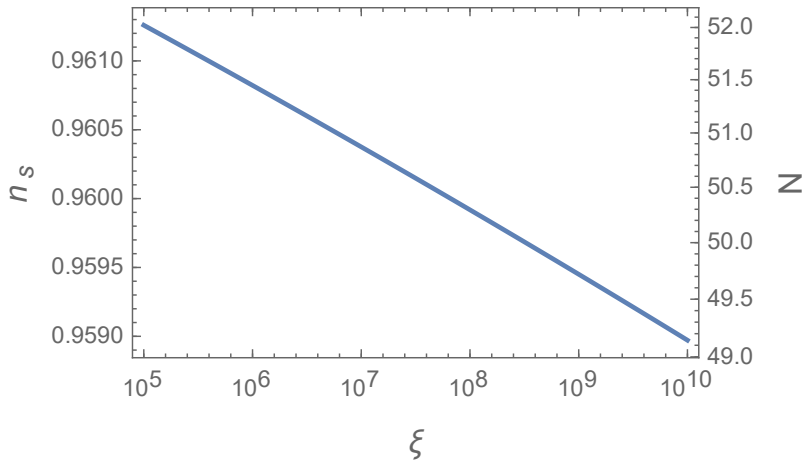
$$N_* = 54.9 - \frac{1}{4} \log \xi$$

Significance on cosmology

- ▶ Spectral index n_s : add 2nd order SR correction and a correction related to SR approximation of N to get

$$n_s \approx 1 - \frac{2}{N_*} - \frac{0.8}{N_*^2}$$

Significance on cosmology



Caveats

- ▶ Non-linear behaviour?
- ▶ Details of backreaction?
- ▶ Thermalization?
- ▶ Quantum corrections to potential ignored

Summary

- ▶ Preheating in Palatini Higgs inflation is very efficient due to tachyonic Higgs production
- ▶ Detailed knowledge of reheating is needed to make accurate cosmological predictions

$$\epsilon_V \approx \frac{1}{8\xi N_V^2}, \quad \eta_V \approx -\frac{1}{N_V}, \quad \xi_V \approx \frac{1}{N_V^2}$$

$$n_s \approx 1 - 6\epsilon_V + 2\eta_V + \frac{1}{3}(44 - 18c)\epsilon_V^2 + (4c - 14)\epsilon_V\eta_V + \frac{2}{3}\eta_V^2 + \frac{1}{6}(13 - 3c)\zeta_V$$

$$N_V \equiv \int_i^x \frac{d\chi}{\sqrt{2\epsilon_V}}$$

$$N_* \approx N_V + 1.8$$

$$n_s \approx 1 - \frac{2}{N_V} + \frac{2.8}{N_V^2} \approx 1 - \frac{2}{N_*} - \frac{0.8}{N_*^2}$$

ξ	Peak			Plateau		
	n_{osc}	ΔN	$\frac{d \ln \rho_{\text{pert}}}{dn_{\text{osc}}}$	n_{osc}	ΔN	$\frac{d \ln \rho_{\text{pert}}}{dn_{\text{osc}}}$
10^5	1.75	0.10	10.3	2	0.07	10.1
10^6	1.25	0.04	12.7	1.5	0.03	12.3
10^7	1.25	0.02	15.0	1.5	0.02	14.2
10^8	0.75	0.007	17.1	1.5	0.009	16.0
10^9	0.75	0.004	19.3	1	0.003	17.9

ξ	$\frac{g_2^2}{\lambda}$	$\langle \frac{\Delta\rho_F}{\rho_B} \rangle$
10^5	27789	0.064
10^6	2814	0.037
10^7	285	0.022
10^8	29	0.014
10^9	2.9	0.014